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TRIGONOMETRICALLY FITTED BLOCK METHOD FOR SOLVING SECOND-ORDER INITIAL VALUE PROBLEMS OF ORDINARY DIFFERENTIAL EQUATIONS

C. O. Alakofa and O. O. Enoch

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RESEARCH ARTICLE

Department of Mathematics, Federal University Oye-Ekiti, Ekiti State, Nigeria

*Corresponding Author E-mail: alakofaoluwaseyi@gmail.com

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ARTICLE DETAILS	ABSTRACT
Article History: Received 02 July 2024 Accepted 05 October 2024 Available online 10 December 2024	This study focuses on developing an implicit continuous four-step method for the numerical solution of initial value problems (IVPs) stemming from ordinary differential equations(ODEs) of second-order. A combination of polynomial and trigonometric basis functions serves as the backbone of this method. The derivation process employs interpolation and collocation techniques, facilitating a robust framework for solving IVPs. Through rigorous analysis employing pertinent theorems, the scheme's consistency, convergence and stability were meticulously scrutinized. The findings of this investigation assert that the developed method exhibits consistency, zero-stability and consequently convergence, underlining its robustness and reliability in practical applications in comparison to existing literature in terms of accuracy, the proposed method outperforms the existing authors when the error are compared. In essence, this work represents a significant advancement in numerical methods for solving IVPs of second-order ODEs. The study provides a strong basis for future research in this area by utilizing a unique combination of polynomial with the sum of sine and cosine as the basis function and putting the approach through theoretical scrutiny. This opens up intriguing directions for further investigation and improvement.
	KEYWORDS

Collocation, Interpolation, Convergence, Stability

Introduction

This paper presents a numerical method for solving initial value problems (IVP) for general second-order ODEs of the form.

$$y^{''} = f(x, y, y^{'}), \ y(a) = \eta_0, \ y^{'}(a) = \eta_1 \ x\varepsilon[a, b]$$

(1)

Second-order linear differential equations find applications in various domains such as physics and engineering, offering a framework to model system of equations. In the realm of numerical integration, particularly for oscillatory problems, the exploration began in 1961 with Gautsci introducing the Adams and St^ormer order using trigonometric polynomials.

Higher order ODEs in ordinary differential equations were first reduced to a system of first-order ODEs, allowing for the application of any suitable method for solving first-order ODEs.Butcher



(2018),Fatunla (1991) and Awoyemi (2001) have studied the reduction procedure in great detail.

Building upon Gautsci groundwork, Psihoyios and Simos (2003, 2005) introduced trigonometrically fitted schemes for solving oscillatory problems. These schemes, operating in predictor-corrector mode, employed the Adams-Bashforth method as the predictor and Adams-Moulton as the corrector. However, their implementation posed challenges, requiring substantial human effort and often resulting in diminished accuracy because the predictor and corrector are usually of different orders.

Addressing these limitations, Jator et. al. (2012) delved into multistep collocation methods to devise trigonometrically fitted approaches based on trigonometric polynomials, introducing Numerov type block methods. Ngwane and Jator (2020) expanded on this by developing a block hybrid scheme specifically tailored for integrating oscillatory problems.

To overcome the drawbacks of the predictor-corrector method, researchers adopted the block method, which can begin computations without requiring previously computed values or additional initialization steps. Noteworthy contributors in this domain include Abolarin et. al. (2020), Omar and Kuboye (2018), Olanegan et. al. (2018) and Awoyemi et. al. (2015) among others.

In this study, an approach involving a second derivative trigonometrically fitted method is presented, utilizing the multistep collocation technique. The approximated interpolating function is constructed as a linear combination of polynomials and trigonometric terms. Specifically, our proposed method is designed to precisely integrate initial value problems (IVPs) stated as combinations of the set $\big\{1,x,x^2,x^3,x^4,\sin(\varpi x),\cos(\varpi x)\big\}_{\!\!\!.\ \ {\rm This\ \ basis\ \ function\ \ is}}$ chosen for its analytical simplicity and its ability to offer an enhanced framework for addressing initial value problems characterized by oscillatory solutions. The combination of these two functions will help to produce more accurate and stable results.

Methodology

In developing this method, the combination of polynomials with sine and cosine functions of the form

$$y(x) = \sum_{n=0}^{k} a_n x^n + a_5 \sin(\varpi x) + a_6 \cos(\varpi x)$$
(2)

is considered as an approximate solution to equation (1). Where the coefficients $a \ s$ are arbitrary real constant that must be equally determined, the frequency of the trigonometric function (\$) will be employed to enhance the method's accuracy.

The second derivative of equation (2) is obtained as:

$$y^{''} = f(x, y, y^{'}), \ \ y(a) = \eta_0, \ \ y^{'}(a) = \eta_1 \ \ x\varepsilon[a, b]$$

Collocating equation (3) at $x = x_{n+j}$, j = 0(1)k and interpolating equation (2) at $x = x_{n+j}$, j = 0,1 gives a seven non-singular equations which can be written as a system in matrix form in equation (4)

l	x_n	x_n^2	x_n^3	x_n^4	$\sin \varpi x_n$	$\cos \varpi x_n$	<i>a</i> ₀		y_n
L	x_{n+1}	x_{n+1}^2	x_{n+1}^3	x_{n+1}^4	$\sin \varpi x_{n+1}$	$\cos \varpi x_{n+1}$	<i>a</i> ₁		y_{n+1}
0	0	2	$6x_n$	$12x_n^2$	$-\varpi^2 \sin \varpi x_n$	$-\varpi^2 \cos \varpi x_n$	a_2		f_n
0	0	2	$6x_{n+1}$	$12x_{n+1}^2$	$-\varpi^2\sin \varpi x_{n+1}$	$-\varpi^2 \cos \varpi x_{n+1}$	a_3	=	f_{n+1}
0	0	2	$6x_{n+2}$	$12x_{n+2}^2$	$-\varpi^2\sin \varpi x_{n+2}$	$-\varpi^2 \cos \varpi x_{n+2}$	a_4		f_{n+2}
0	0	2	$6x_{n+3}$	$12x_{n+3}^2$	$-\varpi^2 \sin \varpi x_{n+3}$	$-\varpi^2 \cos \varpi x_{n+3}$	a_5		f_{n+3}
0	0	2	$6x_{n+4}$	$12x_{n+4}^2$	$-\varpi^2\sin \varpi x_{n+4}$	$-\varpi^2 \cos \varpi x_{n+4}$	a6		f_{n+4}

Then solving equation (4) using Gaussian elimination method with the help of Computer aided software gives values for the unknown parameters $a_n(s)$ which gives our continuous implicit method when substituted into equation (2)

$$y(x) = \sum_{j=0}^{k} \alpha_j(x) y_{n+j} + h^2 \left[\sum_{j=0}^{k} \beta_j(x) f_{n+j} \right]$$
(5)

where $\alpha_j(x)$ and $\beta_j(x)$ are continuous coefficients, $y_{n+j} = y(x_n + jh)$ is the numerical approximation of the analytical solution at x_{n+j} and $f_{n+j} = f(x_{n+j}, y_{n+j}, y'_{n+j})$.

Using $t = \frac{x-x_n}{h}, \frac{dt}{dx} = \frac{1}{h}$, equation (5) is evaluated at the non-interpolating points $x_{n+j} = 2(1)k$ while the

first derivative of Equation (5) is evaluated at all points, their coefficients and the corresponding Taylor series conversion up to $O(u^8)$ by letting u = \$h are given as follows

```
\begin{aligned} & \varphi_{1}(t) = t \\ & \varphi_{1}(t) = t \\ & \varphi_{1}(t) = \frac{1}{2t} \left( \frac{24 \sin(h^{2}) - 22 \sin(h^{2}) + 27 \sin(h^{2}) - 94 \sin(h^{2}) - 96 \sin(h^
```

(3)

(7)

The result is expressed in series form as follows when evaluated at t = 2:

$$\begin{split} y_{n+2} - 2y_{n+1} + y_n &= h^2 \left[\beta_0(x) f_n + \beta_1(x) f_{n+1} + \beta_2(x) f_{n+2} + \beta_3(x) f_{n+3} + \beta_4(x) f_{n+4} \right], \\ \beta_0 &= \frac{19}{240} h^2 + \frac{221}{60480} h^2 u^2 \frac{233}{1814400} h^2 u^4 + \frac{199}{53222400} h^2 u^6 + O(u^8), \\ \beta_1 &= -\frac{17}{20} h^2 + \frac{79}{7560} h^2 u^2 + \frac{79}{226800} h^2 u^4 + \frac{61}{6652800} h^2 u^6 + O(u^8), \\ \beta_2 &= \frac{7}{120} h^2 + \frac{19}{2016} h^2 u^2 + \frac{83}{302400} h^2 u^4 + \frac{1}{197120} h^2 u^6 + O(u^8), \\ \beta_3 &= \frac{1}{60} h^2 - \frac{2}{945} h^2 u^2 - \frac{1}{56700} h^2 u^4 + \frac{1}{415800} h^2 u^6 + O(u^8), \\ \beta_4 &= -\frac{1}{240} h^2 - \frac{31}{60480} h^2 u^2 - \frac{67}{1814400} h^2 u^4 - \frac{109}{53222400} h^2 u^6 + O(u^8) \end{split}$$

The result is expressed in series form as follows when evaluated at t = 3:

$$\begin{split} y_{n+3} - 3y_{n+1} + 2y_n &= h^2 \left[\beta_0(x) f_n + \beta_1(x) f_{n+1} + \beta_2(x) f_{n+2} + \beta_3(x) f_{n+3} + \beta_4(x) f_{n+4} \right], \\ \beta_0 &= \frac{37}{240} h^2 + \frac{137}{20160} h^2 u^2 + \frac{19}{86400} h^2 u^4 \frac{289}{53222400} h^2 u^6 + O(u^8), \\ \beta_1 &= -\frac{9}{5} h^2 + \frac{19}{1008} h^2 u^2 + \frac{83}{151200} h^2 u^4 + \frac{1}{98560} h^2 u^6 + O(u^8), \\ \beta_2 &= \frac{37}{40} h^2 + \frac{53}{3360} h^2 u^2 + \frac{11}{33600} h^2 u^4 - \frac{19}{8870400} h^2 u^6 + O(u^8), \\ \beta_3 &= \frac{2}{15} h^2 - \frac{11}{540} h^2 u^2 + \frac{17}{151200} h^2 u^4 + \frac{173}{13305600} h^2 u^6 + O(u^8), \\ \beta_4 &= \frac{1}{80} h^2 + \frac{31}{20160} h^2 u^2 \frac{67}{604800} h^2 u^4 + \frac{109}{17740800} h^2 u^6 + O(u^8) \end{split}$$

The result is expressed in series form as follows when evaluated at t = 4:

$$\begin{split} y_{n+4} - 4y_{n+1} + 3y_n &= h^2 \left[\beta_0(x) f_n + \beta_1(x) f_{n+1} + \beta_2(x) f_{n+2} + \beta_3(x) f_{n+3} + \beta_4(x) f_{n+4} \right], \\ \beta_0 &= \frac{9}{40} h^2 + \frac{19}{2016} h^2 u^2 + \frac{83}{302400} h^2 u^4 + \frac{1}{197120} h^2 u^6 + O(u^8), \\ \beta_1 &= -\frac{83}{30} h^2 + \frac{37}{1260} h^2 u^2 + \frac{29}{37800} h^2 u^4 + \frac{29}{3226400} h^2 u^6 + O(u^8), \\ \beta_2 &= \frac{37}{20} h^2 + \frac{53}{1680} h^2 u^2 + \frac{11}{16800} h^2 u^4 - \frac{19}{4435200} h^2 u^6 + O(u^8), \\ \beta_3 &= \frac{11}{10} h^2 - \frac{4}{315} h^2 u^2 - \frac{17}{9450} h^2 u^4 + \frac{173}{26611200} h^2 u^6 + O(u^8), \\ \beta_4 &= -\frac{7}{120} h^2 + \frac{11}{10080} h^2 u^2 - \frac{17}{302400} h^2 u^4 - \frac{173}{26611200} h^2 u^6 + O(u^8) \end{split}$$

The block methods are derived by evaluating the first derivative of Equation (5) in order to obtain additional equations needed to couple with (7), (8) and (9), which gives:



(10)

Evaluating the first derivative at all points, and the results are expressed in the series form

$$\begin{split} hy_n' - y_{n+1} + y_n &= h^2 \left[\beta_0(x) f_n + \beta_1(x) f_{n+1} + \beta_2(x) f_{n+2} + \beta_3(x) f_{n+3} + \beta_4(x) f_{n+4} \right], \\ \beta_0' &= -\frac{367}{1440} h^2 - \frac{199}{24192} h^2 u^2 - \frac{1543}{362880} h^2 u^4 - \frac{281}{9123840} h^2 u^6 + O(u^8) \\ \beta_1' &= \frac{3}{8} h^2 - \frac{337}{15120} h^2 u^2 - \frac{457}{453600} h^2 u^4 - \frac{289}{4435200} h^2 u^6 + O(u^8) \\ \beta_2' &= -\frac{47}{240} h^2 - \frac{353}{20160} h^2 u^2 - \frac{19}{40320} h^2 u^4 - \frac{569}{53222400} h^2 u^6 + O(u^8) \\ \beta_3' &= -\frac{29}{360} h^2 + \frac{1}{945} h^2 u^2 - \frac{43}{113400} h^2 u^4 - \frac{127}{2494800} h^2 u^6 + O(u^8) \\ \beta_4' &= -\frac{7}{480} h^2 + \frac{289}{120960} h^2 u^2 + \frac{139}{518400} h^2 u^4 + \frac{2899}{106444800} h^2 u^6 + O(u^8) \\ \end{split}$$

(11)

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$$\begin{split} hy_{n+1}' - y_{n+1} + y_n &= h^2 \left[\beta_0(x) f_n + \beta_1(x) f_{n+1} + \beta_2(x) f_{n+2} + \beta_3(x) f_{n+3} + \beta_4(x) f_{n+4} \right], \\ \beta_0' &= \frac{3}{32} h^2 + \frac{731}{120960} h^2 u^2 + \frac{1439}{3628800} h^2 u^4 + \frac{157}{5068800} h^2 u^6 + O(u^8) \\ \beta_1' &= -\frac{47}{90} h^2 + \frac{97}{6048} h^2 u^2 + \frac{853}{907200} h^2 u^4 + \frac{5317}{79833600} h^2 u^6 + O(u^8) \\ \beta_2' &= -\frac{41}{240} h^2 + \frac{239}{20160} h^2 u^2 + \frac{89}{201600} h^2 u^4 + \frac{743}{53222400} h^2 u^6 + O(u^8) \\ \beta_3' &= \frac{1}{15} h^2 + \frac{1}{4320} h^2 u^2 + \frac{319}{907200} h^2 u^4 + \frac{1277}{26611200} h^2 u^6 + O(u^8) \\ \beta_4' &= -\frac{17}{1440} h^2 - \frac{253}{120960} h^2 u^2 - \frac{181}{725760} h^2 u^4 - \frac{1681}{63866880} h^2 u^6 + O(u^8) \\ \end{split}$$

(8)

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(12)

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$$\begin{split} hy_{n+2}' - y_{n+1} + y_n &= h^2 \left[\beta_0(x) f_n + \beta_1(x) f_{n+1} + \beta_2(x) f_{n+2} + \beta_3(x) f_{n+3} + \beta_4(x) f_{n+4} \right], \\ \beta_0' &= \frac{90}{1440} h^2 + \frac{1}{640} h^2 u^2 - \frac{439}{3628800} h^2 u^4 - \frac{7211}{319334400} h^2 u^6 + O(u^8) \\ \beta_1' &= -\frac{361}{360} h^2 + \frac{29}{5040} h^2 u^2 - \frac{17}{90720} h^2 u^4 - \frac{1817}{39916800} h^2 u^6 + O(u^8) \\ \beta_2' &= \frac{37}{80} h^2 + \frac{53}{6720} h^2 u^2 + \frac{11}{67200} h^2 u^4 - \frac{119}{17740800} h^2 u^6 + O(u^8) \\ \beta_3' &= -\frac{13}{360} h^2 - \frac{1}{210} h^2 u^2 - \frac{23}{56700} h^2 u^4 - \frac{1}{22680} h^2 u^6 + O(u^8) \\ \beta_4' &= -\frac{1}{288} h^2 + \frac{43}{40320} h^2 u^2 + \frac{91}{518400} h^2 u^4 + \frac{7097}{319334400} h^2 u^6 + O(u^8) \\ \end{split}$$

(13)

$$\begin{split} hy_{(n+3)}' - y_{n+1} + y_n &= h^2 [\beta_0(x)f_n + \beta_1(x)f_{n+1} + \beta_2(x)f_{n+2} + \beta_3(x)f_{n+3} + \beta_4(x)f_{n+4}], \\ \beta_0' &= -\frac{119}{1440}h^2 - \frac{571}{120960}h^2u^2 - \frac{1103}{3628800}h^2u^4 - \frac{8291}{319334400}h^2u^6 + O(u^8) \\ \beta_1' &= -\frac{9}{10}h^2 + \frac{65}{6048}h^2u^2 + \frac{517}{907200}h^2u^4 + \frac{59}{1267200}h^2u^6 + O(u^8) \\ \beta_2' &= \frac{263}{240}h^2 + \frac{79}{20160}h^2u^2 - \frac{23}{201600}h^2u^4 - \frac{857}{53222400}h^2u^6 + O(u^8) \\ \beta_3' &= -\frac{4}{9}h^2 + \frac{167}{30240}h^2u^2 + \frac{131}{181440}h^2u^4 + \frac{5431}{79833600}h^2u^6 + O(u^8) \\ \beta_4' &= -\frac{11}{480}h^2 - \frac{59}{17280}h^2u^2 - \frac{1241}{3628800}h^2u^4 - \frac{667}{21288960}h^2u^6 + O(u^8) \end{split}$$

 $hy_{(n+4)}' - y_{n+1} + y_n = -h^2 \left[\beta_0(x)f_n + \beta_1(x)f_{n+1} + \beta_2(x)f_{n+2} + \beta_3(x)f_{n+3} + \beta_4(x)f_{n+4}\right],$

$$\begin{split} \beta_0' &= \quad \frac{9}{160}h^2 + \frac{29}{120960}h^2u^2 - \frac{31}{145152}h^2u^4 - \frac{89}{3225600}h^2u^6 + O(u^8) \\ \beta_1' &= \quad -\frac{377}{360}h^2 + \frac{5}{432}h^2u^2 - \frac{73}{453600}h^2u^4 - \frac{2089}{39916800}h^2u^6 + O(u^8) \\ \beta_2' &= \quad \frac{35}{48}h^2 + \frac{671}{20160}h^2u^2 + \frac{23}{28800}h^2u^4 + \frac{13}{1520640}h^2u^6 + O(u^8) \\ \beta_3' &= \quad \frac{161}{120}h^2 - \frac{31}{945}h^2u^2 - \frac{139}{113400}h^2u^4 - \frac{53}{831600}h^2u^6 + O(u^8) \\ \beta_4' &= \quad \frac{469}{1440}h^2 + \frac{1313}{120960}h^2u^2 + \frac{1741}{3628800}h^2u^4 + \frac{9721}{319334400}h^2u^6 + O(u^8) \end{split}$$

Each of the coefficients in equations (7)-(9) is in trigonometric form. To avoid heavy cancellation that may occur as $u \rightarrow 0$.

Basic Properties of the Block Method

Order and Error Constant

The technique used by Enoch and Alakofa (2024) in finding the order of a method is also adopted in establishing the order of this new developed block method. With this, the proposed block method is of order p = 5 and error constants given by the vector

$$C_7 = \left[\frac{1}{240}, \frac{1}{120}, \frac{1}{120}, \frac{-191}{720}\right]^T$$

Zero Stability of the Block

Definition: A block method is said to be zero stable if as $h \rightarrow 0$, the roots $r_{jj} f = 1(1)k$ of the first characteristic polynomials $\rho(r) = 0$ that is

$$\rho(r) = det\left[\sum A^{(0)}R^{k-1}\right] = 0$$

satisfying $|R| \le 1$, must be simple. $\rho(z) = det[zA^{(0)} - A^{(i)} = 0]$

$$\begin{bmatrix} z \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} - \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix} \end{bmatrix} = 0$$
$$\begin{bmatrix} z & 0 & 0 & -1 \\ 0 & z & 0 & -1 \\ 0 & 0 & z & -1 \\ 0 & 0 & 0 & z -1 \end{bmatrix} = 0$$

since $| z = 1,0,0,0 | \le 1$. Therefore, the block method is zero stable

Consistency

Our new block method is consistent because its order is greater than 1.

Convergence

(14)

(15)

Theorem 3.1: According to Lambert (1973), a linear multistep method is convergent if and only if it is both consistent and zero stable. Thus our block method is convergent since it is zero stable and consistent.

Numerical Experiments

This section examines the performance of the new method by applying it to some second order system of equations. The test problems outcomes are displayed in tabular form.

Error= |exact solution - computed solution|

Test Problem 1

$$y_1'' = \frac{-y_1}{\sqrt{y_1^2 + y_{2^2}}}, y_1(0) = 1, y_1'(0) = 0, h = 0.01$$

$$y_2'' = \frac{-y_2}{\sqrt{y_1^2 + y_{2^2}}}, y_2(0) = 1, y_2'(0) = 1,$$

Theoretical Solution: $y_1 = \cos x$, $y_2 = \sin x$

source: Abdelrahim, R. & Omar, Z. (2016).

	Exact solution of y ₁	Computed solution of <i>y</i> ,	Error in y	Error in Abdelrahim, R. & Omar, Z. (2016)
0.1	0.995004165278 025710	0.99500416542 4475890	1.464501853E (-10)	
0.2	0.980066577841 241630	0.98006657826 3853460	4.226118344E (-10)	4.208611E (-11)

0.3	0.955336489125 605980	0.95533648962 9699970	5.040939888E (-10)	-
0.4	0.921060994002 885100	0.92106099466 0191420	6.573063205E (-10)	2.945689E (-10)
0.5	0.877582561890 372760	0.87758256262 0329190	7.299564286E (-10)	-
0.6	0.825335614909 678330	0.82533561550 3742240	5.940639092E (-10)	7.596528E (-10)
0.7	0.764842187284 488380	0.76484218774 0586540	4.560981592 E(-10)	-
0.8	0.696706709347 165390	0.69670670952 9967600	1.828022178E (-10)	1.178604E (-9)
0.9	0.621609968270 664390	0.62160996813 4783640	1.358807511E (-10)	-
1.0	0.54030230586813 9770	0.540302305363 824840	5.043149232E (-10)	1.180295E (-9)

Table 2: Exact and computed solutions of the new method for solving y_2 inTest Problem 1

		1		
X	Exact solution of y ₂	Computed solution of <i>y</i> ²	Error in y ₂	Error in Abdelrahim, R. & Omar, Z. (2016)
0.1	0.099833416646828	158.099833416645325	85 6 .502298286 (-12)	Е
0.2	0.198669330795061	22 0 .198669330823742	91 0 .868169591 (-11)	E 3.186932E (-10)
0.3	0.295520206661339	\$5 0 .295520206706007 [,]	93 0 .466838011 (-11)	E -
0.4	0.389418342308650	52 0 .389418342373891 [.]	95 6 .524142337 (-11)	E 1.082778E (-9)
0.5	0.479425538604203	01 0 .479425538757451	₽7 0 .532484695 (-10)	E -
0.6	0.564642473395035	37 0 .564642473646759	\$5 0 .517239750 (-10)	E 1.773150E (-9)
0.7	0.644217687237691	130.644217687549621	83 0 .119307035 (-10)	E -
0.8	0.717356090899522	790.717356091224827	24 8 .253044500 (-10)	E 1.862680E (-9)
0.9	0.783326909627483	410.783326909947776	32 0 .202929033 (-10)	E -
1.0	0.841470984807896	\$0 0 .841470985119705	19 8 .118086900 (-10)	E 1.081632E (-9)

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Test Problem 2: (Periodic Problem) Vande Vyver

0.

0.

$$y_1'' + y =$$

$$y_2'' + y =$$

$$001\sin(x), y_2(0) = 0, y_2'(0) = 0.999$$

 $001\cos(x), y_1(0) = 1, y'_1(0) = 0, h =$

Theoretical Solution:

 $y_1(x) = \cos(x) + 0.0005x\sin(x),$

 $y_2(x) = \sin(x) - 0.0005x\cos(x).$

source: Kayode et. al (2021)

Exact

0.99999995005

0.9999980020

0.9999955045

0.9999920080

0.9999875125

9999820180

×

0.001

0.002

0.003

0.004

0.005

0.006

0.007	ъ	0.999975524599895	5.306866058E	2.44963093E
	0.9999755245 99842		(-14)	(-05)
0.008	1	0.999968032170401	7.593925488E	3.20040361E
	0.9999680321 70325		(-14)	(-05)
0.009	7	0.999959540772948	1.210143097E	4.05095821E
	0.9999595407 72827		(-13)	(-05)
0.01	4	0.999950050416012	1.800781746E	5.00240831E
	0.9999500504 15832		(-13)	(-05)

Table 4: Exact and computed solutions of the new method forsolving y2 in Test Problem 2

Error in y₂

(-06)

(-06)

(-05)

(-05)

(-05)

(-05)

4.499999217E

8.999994677E

1.349998254E

1.799995830E

2.249991885E

2.699985866E

in Kayode et. al (2021)

4.98251761E

1.99750245E

4.49460223E

7.99469593E

1.24926327E

1.79955544E

Error

(-07)

(-06)

(-06)

(-06)

(-05)

(-05)

	Computed	Error in <i>y</i> 1	Error in Kayode et. al (2021)	x	Exact	Computed
				0.001	0.0009994998 33583	0.000994999834367
	0.999999500500042	0	4.98251761E (-07)		0.00	
00042				0.002	0.0019989986 68667	0.001989998673990
	0.999998002000666	9.992007222E (-16)	1.99750245E (-06)		0.00	
00665				0.003	0.0029984955 06752	0.002984995524214
	0.999995504503371	2.886579864E (-15)	4.49460223E (-06)		0.00	
03368				0.004	0.0039979893 49342	0.003979989391042
	0.999992008010650	5.107025913E (-15)	7.99469593E (-06)		0.00	
10645				0.005	0.0049974791 97943	0.004974979279097
0	0.999987512526005	1.576516695E (-14)	1.24926327E (-05)		0.00	
25990				0.006	0.0059969640 54065	0.005969964195402
	0.999982018053925	3.264055692E (-14)	1.79955544E (-05)		0.00	
53892		(- 1)	(33)			

Table 3: Exact nd computed solutions of the new method for solving y_1 in Test Problem 2

0				
	hafe and O. O. Freech (2024) <i>Tuis</i> an am	atui anllu. Pitta d	Die de Mathead fan ande in a Cananad Ondar Initial Value Drahlanne of Ordinarra
	Kofa ana U. U. Enoch. (2 Forential Equations	2024). Trigonom	etrically Fittea	Block Method for solving Second-Order Initial Value Problems of Ordinary

0.007	0.0069964429 19223	0.006964943147947	3.149977128E (-05)	2.44963093E (-05)
0.008	0.0079959147 94939	0.007959915142753	3.599965219E (-05)	3.20040361E (-05)
0.009	0.0089953786 82741	0.008954879181453	4.049950129E (-05)	4.05095821E (-05)
0.01	0.0099948335 84165	0.009949834271476	4.499931269E (-05)	5.00240831E (-05)

Discussion of Result and Conclusion

The numerical results produced from the analyzed problems demonstrate high efficiency and effectiveness in their performance. From the results generated, it is worthy of note the favorable performance of the proposed method when compared with existing methods in the literature. Tables 1 and 2 presents the computational solution of test problem 1 (systems of equation), the results compete favorably well with that of Abdelrahim and Omar (2016). in terms of errors. Also, Tables 3 and 4 shows the computational solution of test problem 2 (VandeVyver), the results outperform that of Kayode et. al (2021). Utilizing collocation and interpolation techniques, a new class of continuous second derivative block methods for solving ODEs is constructed. This innovative approach combines polynomial and trigonometric functions, implemented through code written in MATLAB, to develop an approximation solution. The resulting block techniques exhibit continuous coefficients and possess key properties of consistency, zero stability, and convergence. These characteristics contribute to a robust methodology that ensures reliability in solving ODEs, presenting a promising avenue for further exploration and application in this field

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